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RESONANCE SCATTERING OF GAMMA QUANTA BY Mg^{24} NUCLEI

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Resonance scattering of 1.38-MeV γ quanta by Mg^{24} nuclei was investigated using a source of radioactive Na^{24} in the form of an aqueous solution of NaOH. The average cross section for resonance scattering was found to be $(3.7 \pm 0.6) \times 10^{-28}$ cm². The energy distribution of the emitted γ quanta was calculated for the β - γ_2 - γ_1 cascade taking account of the slowing down of the recoil nuclei. The lifetime of the 1.38-MeV level in Mg^{24} was determined to be $(1.1 \pm 0.2) \times 10^{-12}$ sec.

RECENT investigations^[1-3] have shown that resonance scattering of γ quanta by nuclei can occur when one uses liquid and solid sources, provided the lifetime τ_γ of the initial level is less than 10^{-12} sec. Since the time between collisions of the recoil nuclei with surrounding atoms is of order 10^{-13} – 10^{-14} sec, the slowing down of the recoil nuclei results in a change in the energy distribution $P(E)$ of the γ quanta.

To determine τ_γ one must find experimentally the average cross section for resonance scattering and calculate the microspectrum $P(E)$ theoretically. A calculation of the microspectrum, with account of the slowing down of the recoil nuclei, is necessary mainly for two reasons. First, in certain cases, in particular when using very short-lived isotopes, obtaining a source in the gaseous state or the use of the "self-absorption" method for determining τ_γ is a difficult experimental problem. Second, the study of the attenuation of the resonance effect in dense media can give information about the lifetimes of higher-lying levels.

In our work we studied the resonance scattering of 1.38-MeV γ quanta by Mg^{24} nuclei, using the radioactive isotope Na^{24} ($T_{1/2} = 14$ hr) in the form of an aqueous solution of NaOH. The Na^{24} decay is accompanied by the emission of a β particle with endpoint $E_{\max} = 1.39$ MeV and two γ quanta with energies $E_{\gamma_2} = 2.76$ MeV and $E_{\gamma_1} = 1.38$ MeV.

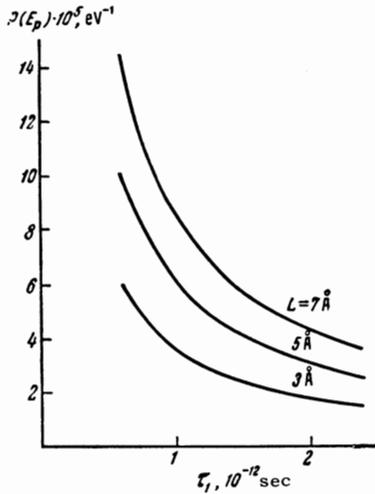
To compute the spectrum $P(E)$, we used a model with continuous slowing down of the recoil nuclei as the result of elastic collisions with the surrounding atoms. The energy of the emitted γ_1 quantum was

$$E_{\gamma_1} = (E_{\gamma_1^0} / Mc) p_{2x} + [E_{\gamma_1^0} - (E_{\gamma_1^0})^2 / 2Mc^2] = ap_{2x} + b,$$

where $E_{\gamma_1^0}$ is the energy of the level, p_{2x} is the projection of the total momentum on the direction of emergence of the γ_1 quantum, and the energy distribution of the γ quanta is given by the density distribution of the momentum projection $\xi(p_{2x}')$:

$$P(E) = \xi(p_{2x}') / a. \quad (1)$$

For the case of a β - γ_2 - γ_1 cascade, the density distribution of recoil nuclei over momentum pro-



Dependence of $P(E_p)$ on τ_1 , computed theoretically for various values of the free path L .

jections, taking account of the slowing down, was found to be

$$\xi(p_{2x''}) = \frac{M^5 L^5}{4 p_{\gamma_2} \tau_1 \tau_2} \int_{|p_{2x''}|}^{p_{\max} + p_{\gamma_2}} dp_{2x''} \int_{\lambda_2}^{\lambda_1} \frac{\exp(-t_1/\tau_1) dt_1}{(ML - p_2'' t)^3} \\ \times \int_0^{p_{\max}} \frac{dp'}{p'} \int_0^{\lambda_0} \exp\left(-\frac{t}{\tau_2}\right) g\left(\frac{MLp'}{ML - p't}\right) \\ \times \frac{dt}{(ML - p't)^2}, \\ \lambda_0 = ML \left(\frac{1}{p'} - \frac{1}{p_{\max}} \right), \quad \lambda_1 = ML \left(\frac{1}{p''} - \frac{1}{p_{\gamma_2} + p_{\max}} \right), \\ \lambda_2 = ML \left(\frac{1}{p''} - \frac{1}{p_{\gamma_2} - p_{\max} - p_2''} \right). \quad (2)$$

Here $g(p)$ is the momentum distribution of the recoil nuclei, p_{\max} is the maximum momentum after the β decay, p_2'' is the total momentum after the β and γ_2 transitions, p_{γ_2} is the momentum of the recoil nucleus from the γ_2 transition, τ_1 and τ_2 are the lifetimes of the first and second excited states, and L is the mean free path.

Theoretical values of $\xi(p_{2x''})$ were computed on the "Minsk" computer for the cases of the $\text{Na}^{24} \rightarrow \text{Mg}^{24}$ and $\text{Co}^{60} \rightarrow \text{Ni}^{60}$ decays, with different values of τ_1 , τ_2 , and L . The computations showed that for $\text{Co}^{60} \rightarrow \text{Ni}^{60}$ the value of $P(E_p)$ depends on τ_2 , which is easily understood considering that $E_{\gamma_2} < E_{\gamma_1}$. For $\text{Na}^{24} \rightarrow \text{Mg}^{24}$ ($E_{\gamma_2} > E_{\gamma_1}$), the value of $P(E_p)$ depends only on τ_1 and L and is independent of τ_2 , which is explained by the position of the resonance line in the microspectrum.

The figure shows the dependence found for $P(E_p)$ on τ_1 for different values of L . Consequently, assigning L , we can determine τ_1 if we know the average cross section $\bar{\sigma}$ for resonance scattering of 1.38-MeV γ quanta by Mg^{24} nuclei. The determination of $\bar{\sigma}$ was done in an apparatus similar to one used earlier.^[4] The initial Na^{24} activity was 0.5 Ci. Since the resonance effect also occurs in solid Na^{24} ,^[1,5] to measure the scattering from scatterers of magnesium and aluminium under nonresonant conditions we used the radioactive isotope La^{140} ($T_{1/2} = 40$ hr) with transition energy 1.59 MeV. The determination of the average cross section for resonance scattering was done by numerical integration over the scatterer surface. We found the value $\bar{\sigma} = (3.7 \pm 0.6) \times 10^{-28} \text{ cm}^2$.

Analysis of the data found from resonance scattering of γ quanta emitted after a β transition^[4,3] shows that the mean free path of the recoil nuclei in water solution is $(1.5 \pm 0.3) \times 10^{-8} \text{ cm}$, which gives $L = (5.0 \pm 1.0) \times 10^{-8} \text{ cm}$ for the case considered. Using the experimental value of $\bar{\sigma}$ and the theoretical dependence of $P(E_p)$ on τ_1 , we find

$$\tau_1 = (1.1 \pm 0.2) \cdot 10^{-12} \text{ sec.}$$

The result is in satisfactory agreement with the data of other work on resonance scattering of 1.38-MeV γ quanta by Mg^{24} nuclei.

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