

A FORM FACTOR SATISFYING THE UNITARITY CONDITION

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Submitted to JETP editor May 22, 1962

 J. Exptl. Theoret. Phys. (U.S.S.R.) **43**, 1769-1770 (November, 1962)

A simple relation between the scattering amplitude and particle form factor is established on the basis of the unitarity condition by applying the R-matrix formalism.

A form factor satisfying the unitarity and analyticity conditions is usually obtained in terms of the scattering phase shift as a solution of a singular integral equation^[1,2]. It is somewhat easier to obtain the solution by using the R-matrix formalism, where under simple assumptions concerning the character of the scattering there is no need to solve the integral equations.

The two-particle unitarity condition for the partial amplitude $T(s)$

$$T^* - T = 2i\rho T^*T,$$

$$\rho = s^{-1} \{ [s - (m_1 + m_2)^2] [s - (m_1 - m_2)^2] \}^{1/2}$$

(the factor ρ corresponds to the statistical weight of the two-particle state) is satisfied identically if we put

$$T = R(1 - i\rho R)^{-1} \equiv (\rho'/\rho) R' (1 - i\rho' R')^{-1},$$

$$\rho' = \sqrt{s - (m_1 + m_2)^2}, \quad (1)$$

where R' does not have a right-hand cut but has all other singularities of T . Such a simple solution is applicable because the two-particle cut corresponds to an ambiguity which is guaranteed by the fact that ρ' is in the form of a square root.

This fact uncovers the possibility of a "semi-phenomenological" description of scattering in a certain limited region of values of s by suitable choice of R' ^[3,4]. Putting $R' = \text{const}$, we obtain an approximation of the scattering length. If R' is a polynomial in s , this corresponds to expansion in powers of the "effective radius." The pole R' on the real axis describes a Breit-Wigner type of resonance.

The unitarity condition for the form factor $F(s)$ (here s is square of the momentum transfer) is, if only one partial wave is taken into account,

$$F^* - F = 2i\rho F^*T. \quad (2)$$

Substituting (1) in (2) we verify that $F = f(1 - i\rho'R')^{-1}$ satisfies (2) identically for any $f(s)$ which has no right-hand cut. It is necessary to choose f such as to obtain the correct analytic properties of $F(s)$.

$1 - i\rho'R'$ can have a left-hand cut corresponding to the cut of $T(s)$. The function f should then contain a left-hand cut, and its discontinuity is chosen such as to make the form factor F an analytic function in the left-plane. We thus obtain for f the Muskhelishvili-Omnès equation

$$\frac{\text{Im} f}{f} = -\frac{|\beta'| |\text{Im} R'|}{1 + |\rho'| R'}, \quad (3)$$

which, as is well known, is solved with accuracy to an arbitrary polynomial, the degree of which determines the behavior of f when $s \rightarrow \infty$. The function f also has poles coinciding with the poles of R' . The requirement that F have no poles means that f have no zeros and have all the poles of R' . Then f is determined uniquely as the fundamental solution of (3) if R' is known, i.e., if the scattering amplitude is known.

In a small region of variation of s the left-hand cut can be approximated by a finite number of constants, and R' can be assumed to have no cuts, i.e., to be a rational function of s . Then f is also a rational function of s , which has no other poles except those possessed by R' . This function is determined completely by the zeros of F . If we do not specify the zeros of F , then to reduce the arbitrariness in the choice of f we can use the requirement, reasonable under these conditions, that $|F|$ not increase when $s \rightarrow \infty$, a requirement which determines the maximum degree of f .

Let us consider by way of an example the behavior of F near the values $s = s_0$ at which resonance takes place in the scattering:

$$R' = \frac{a}{s - s_0}, \quad F = \frac{f(s - s_0)}{s - s_0 - i\rho'a}.$$

The requirements that $|F|$ be non-increasing and $F(0) = 1$ yield

$$f \sim \frac{s - s_1}{s - s_0}, \quad F = \frac{s_0 - (m_1 + m_2)a}{s_1} \frac{s - s_1}{s - s_0 - i\rho'a}.$$

If $F \neq 0$ in the region of interest to us, i.e., the zero of F (the point s_1) is situated far away from s_0 , then we can assume that F has a zero

at infinity. Then, obviously,

$$F = \frac{-s_0 + (m_1 + m_2) a}{s - s_0 - i\rho a}.$$

Grashin and Mel'nikov^[2,5] have pointed out that the zero s_1 of the form factor can be determined uniquely from the condition that the inelastic processes have a small influence, with s_1 situated in the resonant region. However, without insisting on the point of view that analyticity and unitarity must determine uniquely the scattering-matrix elements, such a choice of the solution turns out to be arbitrary. The requirement that an infinitesimally narrow resonance in the scattering make no contribution to the form factor, i.e., that $s_1 \rightarrow s_0$ as $a \rightarrow 0$, does not contradict the example presented, in which $|s_1 - s_0|/ma \gg 1$ for finite a .

Thus, in a small region of momentum-transfer values it is possible to obtain in very simple fashion a form factor satisfying the unitarity condition in the two-particle approximation. At the same time, the ambiguity typical of problems of this

kind, first pointed out by Castillejo, Dalitz, and Dyson, arises in this case.

I am glad of the opportunity to express my gratitude to B. L. Ioffe and L. P. Rudik for suggesting the problem and for valuable advice.

¹R. Omnes, *Nuovo cimento* 8, 316 (1958).

²A. F. Grashin, *JETP* 43, 277 (1962), *Soviet Phys. JETP* 16, 198 (1963).

³R. H. Dalitz and S. F. Tuan, *Ann. of Phys.* 10, 307 (1960).

⁴B. L. Ioffe, in *Coll. Voprosy teorii vzaimodei-stviya elementarnykh chastits* (Problems in the Theory of Elementary Particle Interaction), Erevan, in press.

⁵A. F. Grashin and V. N. Mel'nikov, *JETP* 42, 1404 (1962), *Soviet Phys. JETP* 15, 973 (1962).

⁶Castillejo, Dalitz, and Dyson, *Phys. Rev.* 101, 453 (1956).

Translated by J. G. Adashko.
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