

# Letters to the Editor

## CONCERNING THE PAPER "EVALUATION OF PHASE INTEGRALS IN THE COVARIANT FORMULATION OF THE THEORY OF MULTIPLE PRODUCTION OF PARTICLES" BY L. G. YAKOVLEV

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THE paper by Yakovlev<sup>1</sup> contains basically a proposal to utilize the simple integral (12), over the multidimensional region defined by Eq. (16) and Fig. 1, for the phase volume evaluation. We shall show that the integral (12) is incorrect and that it is impossible to regard (16) as representing the limits of integration.

Deriving Eq. (12) from (11), the author integrates over the directions of the particle momenta after fixing their magnitudes. While doing so, he assumes the possibility of having arbitrary angles between the momentum which is being integrated and the vector sum of those not yet integrated. However, when one of the momenta is close enough to the limiting value, the angles mentioned above cannot differ very much from 180°. In the case of similar (and many other) states there exists a limiting angle which is a function of the chosen momentum values. In the presence of such a limiting angle [cf. reference 2, Eq. (17)] the limits of integration become variable, and one cannot derive Eq. (12).

Further, it is easy to show that the portions of the region of integration (16) close to the vertices of the hexagon in Fig. 1 of reference 1 do not correspond to physical states of a system consisting of three zero-rest-mass particles. In the region close to the right lower vertex, particle 3 should come almost to rest and simultaneously move at almost the limiting velocity. It is equally true that when  $n > 3$  (16) does not represent the region of integration but a polyhedron circumscribed around it. The author errs here by accepting the necessary conditions imposed on the energy (16) as necessary and sufficient. If the sufficient conditions are also taken into account, the region of integration is not bounded by planes but by some very complicated curved surfaces (plotted roughly in Fig. 1 of reference 3 and given analytically by

Eq. (2.13) of the same paper). The character of these surfaces does not permit simple integrations as those attempted by Yakovlev.

It follows from the above that the simplicity of the method for evaluation of covariant weights was achieved by means of an incorrect extension of the region of integration to nonphysical states, i.e., to impossible values of angles and energy. Inaccurate are, in particular, Eqs. (12), (13), and (17) to (19), and the estimate of the applicability of the correct formula (9) is questionable.

<sup>1</sup> L. G. Yakovlev, JETP **37**, 1041 (1959), Soviet Phys. JETP **10**, 741 (1960).

<sup>2</sup> Yu. N. Blagoveshchenskiĭ and G. I. Kopylov, Preprint R-213, Joint Institute for Nuclear Research (U.S.S.R.) (1958).

<sup>3</sup> G. I. Kopylov, JETP **35**, 1426 (1958), Soviet Phys. JETP **8**, 996 (1959).

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## SIGNS OF CONSTANTS OF STRONG INTERACTIONS

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STRONG interactions of known elementary particles are characterized by eight vertices

$$NN\pi, \Lambda\Sigma\pi, \Sigma\Sigma\pi, \Xi\Xi\pi, N\Lambda K, N\Sigma K, \Xi\Lambda K, \Xi\Sigma K, \quad (1)$$

to each of which corresponds a respective coupling constant. The only one known for the present is the  $NN\pi$ -interaction constant. The determination of the remaining seven constants represents one of the basic tasks of high energy physics. A fundamental characteristic of a vertex is not only the absolute magnitude of the coupling constant but also its sign. The knowledge of the signs of the constants is of special importance in inquiries into various symmetry properties characterizing strong interactions.

It is obvious that the absolute sign of a particular constant does not have physical meaning, since it can always be associated with the field of one of the particles entering the vertex. Only a product of signs of vertices whose aggregate contains any